DEEPLY VIRTUAL COMPTON SCATTERING WITH CLAS6

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JLab12 - Italy meeting, 13.07.2015



Introduction

GPDs describe the nucleon structure in terms of quark and gluon degrees of freedom

Correlation between transverse position and longitudinal momentum fraction of quark in the nucleon



Form Factors: Transverse distribution of quarks in space coordinate. $F(t)=\int dx^*GPD(x, \xi, t)$ PDFs: Quark longitudinal momentum fraction in the nucleon. $q(x)=GPD(x, \xi=0, t=0)$

Relation to total angular momentum (Ji relation):

$$J_q = \frac{1}{2} \lim_{t \to 0} \int_{-1}^1 dx \, x \left(H_q(x,\xi,t) + E_q(x,\xi,t) \right)$$

leading-twist, quark chirality conserving		
spin-1/2	unpolarized	polarized
no nucleon hel. flip	H	\widetilde{H}
nucleon hel. flip	E	\widetilde{E}



Experimental Access to GPDs

Experimental probe of GPDs — Hard exclusive Processes



Deeply Virtual Compton Scattering

- Theoretically the cleanest probe of GPDs
- Theoretical accuracy at NNLO
- GPDs are accessed through convolution integrals with hard scattering amplitude
- Experimental observables: Azimuthal asymmetries, cross sections, cross section differences. \sim
- Amplitudes depend on all GPDs H, E, H, E

DVCS and Bethe-Heitler \Rightarrow Same final state \Rightarrow Interference $\frac{d\sigma}{dx_B dQ^2 d|t| d\phi} \propto |\mathcal{T}_{BH}|^2 + |\mathcal{T}_{DVCS}|^2 + \underbrace{\mathcal{T}_{DVCS} \mathcal{T}^*_{BH} + \mathcal{T}_{BH} \mathcal{T}^*_{DVCS}}_{I}$

- At CLAS kinematics $|\mathcal{T}_{DVCS}|^2 << |\mathcal{T}_{BH}|^2$
- DVCS amplitudes can be accessed trough Interference
- Interference \Rightarrow non-zero azimuthal asymmetries



Deeply Virtual Compton Scattering





 $\int \Phi^{*} \Phi$

 $(\overrightarrow{P}, \overrightarrow{e}, \overrightarrow{\phi}) = \mathcal{O}(\overrightarrow{P}, \overrightarrow{e}, \phi) \propto \operatorname{Re}[\overrightarrow{F_1}, \overrightarrow{H}] (\operatorname{Re}[\overrightarrow{G_1}, \overrightarrow{H_1}]) \operatorname{GPD's}$ $\mathcal{F}(\xi, t) = \sum_{q} \int_{-1}^{1} dx C_q(\xi, x) F^q(x, \xi, t)$



Azimuthal Dependence



DVCS at **CLAS**

DVCS with Longitudinally Polarized Target :

Analysis based on egl-dvcs experiment

Extraction of beam-spin, target-spin and double-spin asymmetries

 $d\sigma = d\sigma_{UU} + P_{Beam} d\sigma_{LU} + P_{Target} d\sigma_{UL} + P_{Beam} P_{Target} d\sigma_{LL}$ **DVCS channel selection**

Detection of complete final state of the process $\,ep \to ep\gamma\,$





DVCS at **CLAS**

Single Beam-Spin & Target-Spin Asymmetries

S.Pisano et al. Phys.Rev.D91:052014 (2015)





DVCS at **CLAS**

Longitudinal Double-Spin Asymmetry

S.Pisano et al. Phys.Rev.D91:052014 (2015)







Single Beam-Spin & Target-Spin Asymmetries

☞ Kresimir Kumericki: Status of GPD phenomenology - fits to DVCS data (IWHSS 2015, Suzdal, Russia)





Longitudinal Double-Spin Asymmetry

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Cross Section Results

4-fold differential cross-sections & cross-section difference



$$\begin{aligned} |\mathcal{T}_{\rm BH}|^2 &= \frac{K_{\rm BH}}{\mathcal{P}_1(\phi)\mathcal{P}_2(\phi)} \Biggl\{ \sum_{n=0}^2 c_n^{\rm BH} \cos(n\phi) \Biggr\} \\ |\mathcal{T}_{\rm DVCS}|^2 &= K_{\rm DVCS} \Biggl\{ \sum_{n=0}^2 c_n^{\rm DVCS} \cos(n\phi) + \sum_{n=1}^2 s_n^{\rm DVCS} \sin(n\phi) \Biggr\} \\ \mathcal{I} &= -\frac{K_{\rm I}e_\ell}{\mathcal{P}_1(\phi)\mathcal{P}_2(\phi)} \Biggl\{ \sum_{n=0}^3 c_n^{\rm I} \cos(n\phi) + \sum_{n=1}^3 s_n^{\rm I} \sin(n\phi) \Biggr\} \end{aligned}$$

H.S.Jo et al. Phys.Rev.Lett. (2015) [submitted]





Cross Sections & cross section differences

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Ongoing DVCS analysis

DVCS measurement via $ep \rightarrow epX$ Analysis based on el6 data data (possibility to combine with elf)









Ongoing DVCS analysis

Exclusive pion production $ep \rightarrow ep\gamma\gamma$

Data -MC comparison





Ongoing DVCS analysis

DVCS with photon detection $~ep \rightarrow ep\gamma$

Data -MC comparison





Thank you.



Longitudinal Double-Spin Asymmetry

Resimir Kumericki: Status of GPD phenomenology - fits to DVCS data





